#### THE TWO DIMENSIONAL HANNAY-BERRY MODEL

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#### Abstract

The main goal of this paper is to construct the Hannay-Berry model of quantum mechanics, on a two dimensional symplectic torus. We construct a simultaneous quantization of the algebra of functions and the linear symplectic group  $\Gamma = \operatorname{SL}_2(\mathbb{Z})$ . We obtain the quantization via an action of  $\Gamma$  on the set of equivalence classes of irreducible representations of Rieffel's quantum torus  $\mathcal{A}_{\hbar}$ . For  $\hbar \in \mathbb{Q}$  this action has a unique fixed point. This gives a canonical projective equivariant quantization. There exists a Hilbert space on which both  $\Gamma$  and  $\mathcal{A}_{\hbar}$  act equivariantly. Combined with the fact that every projective representation of  $\Gamma$  can be lifted to a linear representation, we also obtain linear equivariant quantization.

# 0 Introduction

### 0.1 Motivation

In the paper "Quantization of linear maps on the torus - Fresnel diffraction by a periodic grating", published in 1980 (cf. [HB]), the physicists J. Hannay and M.V. Berry explore a model for quantum mechanics on the 2-dimensional torus. Hannay and Berry suggested to quantize simultaneously the functions on the torus and the linear symplectic group  $\Gamma = \mathrm{SL}_2(\mathbb{Z})$ . They found (cf. [HB],[Me]) that the theta subgroup  $\Gamma_{\Theta} \subset \Gamma$  is the largest that one can quantize and asked (cf. [HB],[Me]) whether the quantization of  $\Gamma$  satisfy a multiplicativity property (i.e., is a linear representation of the group). In this paper we want to construct the Hannay-Berry's model for the bigger group of symmetries, i.e., the whole symplectic group  $\Gamma$ . The central question is whether there exists a Hilbert space on which a deformation of the algebra of functions and the linear symplectic group  $\Gamma$  both act in a compatible way.

### 0.2 Results

In this paper we give an affirmative answer to the existence of the quantization procedure. We show a construction (Theorem 0.3, Corollary 0.4 and Theorem 0.5) of the canonical equivariant quantization procedure for rational Planck constants. It is unique as a projective quantization (see definitions below). We show that the projective representation of  $\Gamma$  can be lifted in exactly 12 different ways to a linear representation (to obey the multiplicativity property). These are the first examples of such equivariant quantization for the whole symplectic group  $\Gamma$ . Our construction slightly improves the known constructions [HB, Me, KR1] for which the group of quantizable elements is  $\Gamma_{\Theta} \subset \Gamma$  and gives a positive answer to the Hannay-Berry question on the linearization of the projective representation of the group of quantizable elements. (cf. [HB], [Me]). Previously it was shown by Mezzadri and Kurlberg-Rudnick (cf. [Me], [KR1]) that one can construct an equivariant quantization for the theta subgroup, in case when the Planck constant is of the form  $\hbar = \frac{1}{N}$ ,  $N \in \mathbb{N}$ .

#### 0.2.1 Classical torus

Let  $(\mathbf{T}, \omega)$  be the two dimensional symplectic torus. Together with its linear symplectomorphisms  $\Gamma \simeq \mathrm{SL}_2(\mathbb{Z})$  it serves as a simple model of classical mechanics (a compact version of the phase space of the harmonic oscillator). More precisely, let  $\mathbf{T} = \mathrm{W}/\Lambda$  where W is a two dimensional real vector space, i.e.,  $\mathrm{W} \simeq \mathbb{R}^2$  and  $\Lambda$  is a rank two lattice in W, i.e.,  $\Lambda \simeq \mathbb{Z}^2$ . We obtain the symplectic form on  $\mathbf{T}$  by taking a non-degenerate symplectic form on W:

$$\omega: W \times W \longrightarrow \mathbb{R}.$$

We require  $\omega$  to be integral, namely  $\omega : \Lambda \times \Lambda \longrightarrow \mathbb{Z}$  and normalized, i.e.,  $Vol(\mathbf{T}) = 1$ .

Let  $\mathrm{Sp}(\mathrm{W},\omega)$  be the group of linear symplectomorphisms, i.e.,  $\mathrm{Sp}(\mathrm{W},\omega) \simeq \mathrm{SL}_2(\mathbb{R})$ . Consider the subgroup  $\Gamma \subset \mathrm{Sp}(\mathrm{W},\omega)$  of elements that preserve the lattice  $\Lambda$ , i.e.,  $\Gamma(\Lambda) \subseteq \Lambda$ . Then  $\Gamma \simeq \mathrm{SL}_2(\mathbb{Z})$ . The subgroup  $\Gamma$  is the group of linear symplectomorphisms of  $\mathbf{T}$ .

We denote by  $\Lambda^* \subseteq W^*$  the dual lattice:

$$\Lambda^* := \{ \xi \in W^* | \xi(\Lambda) \subset \mathbb{Z} \}.$$

The lattice  $\Lambda^*$  is identified with the lattice  $\mathbf{T}^{\vee} := \operatorname{Hom}(\mathbf{T}, \mathbb{C}^*)$  of characters of  $\mathbf{T}$  by the following map:

$$\xi \in \Lambda^* \longmapsto e^{2\pi i < \xi, \cdot>} \in \mathbf{T}^{\vee}.$$

The form  $\omega$  allows us to identify the vector spaces W and W\*. For simplicity we will denote the induced form on W\* also by  $\omega$ .

#### 0.2.2 Equivariant quantization of the torus

We will construct a particular type of quantization procedure for the functions. Moreover this quantization will be equivariant with respect to the action of the "classical symmetries"  $\Gamma$ :

**Definition 0.1** By Weyl quantization of  $\mathcal{A}$  we mean a family of  $\mathbb{C}$ -linear, \*- morphisms  $\pi_{\hbar}: \mathcal{A} \longrightarrow End(\mathcal{H}_{\hbar}), \ \hbar \in \mathbb{R}$ , where  $\mathcal{H}_{\hbar}$  is a Hilbert space, s.t. the following property holds:

$$\pi_{\scriptscriptstyle\hbar}(\xi+\eta) = e^{\pi i \hbar w(\xi,\eta)} \pi_{\scriptscriptstyle\hbar}(\xi) \pi_{\scriptscriptstyle\hbar}(\eta)$$

for all  $\xi, \eta \in \Lambda^*$  and  $\hbar \in \mathbb{R}$ .

This type of quantization procedure will obey the "usual" properties (cf. [D4]):

$$\begin{split} ||\pi_{{\scriptscriptstyle\hbar}}(fg)-\pi_{{\scriptscriptstyle\hbar}}(f)\pi_{{\scriptscriptstyle\hbar}}(g)||_{_{\mathcal{H}_{\hbar}}} & \longrightarrow & 0, \quad as \; \hbar \to 0, \\ ||\frac{i}{\hbar}[\pi_{{\scriptscriptstyle\hbar}}(f),\pi_{{\scriptscriptstyle\hbar}}(g)]-\pi_{{\scriptscriptstyle\hbar}}(\{f,g\})||_{_{\mathcal{H}_{\hbar}}} & \longrightarrow & 0, \quad as \; \hbar \to 0. \end{split}$$

where  $\{,\}$  is the Poisson brackets on functions.

**Definition 0.2** By equivariant quantization of  $\mathbf{T}$  we mean a quantization of  $\mathcal{A}$  with additional maps  $\rho_h : \Gamma \longrightarrow U(\mathcal{H}_h)$  s.t. the following equivariant property (called Egorov's identity) holds:

$$\rho_{h}(B)^{-1}\pi_{h}(f)\rho_{h}(B) = \pi_{h}(f \circ B)$$
(0.2.1)

for all  $\hbar \in \mathbb{R}$ ,  $f \in \mathcal{A}$  and  $B \in \Gamma$ . Here  $U(\mathcal{H}_{\hbar})$  is the group of unitary operators on  $\mathcal{H}_{\hbar}$ . If  $(\rho_{\hbar}, \mathcal{H}_{\hbar})$  is a projective (respectively linear) representation of the group  $\Gamma$  then we call the quantization projective (respectively linear).

The idea of the construction is as follows: We use a "deformation" of the algebra  $\mathcal{A}$  of functions on  $\mathbf{T}$ . We define an algebra  $\mathcal{A}_{\hbar}$ , usually called the two dimensional non-commutative torus (cf. [Ri]). If  $\hbar = \frac{M}{N} \in \mathbb{Q}$ , then we will see that all irreducible representations of  $\mathcal{A}_{\hbar}$  have dimension N. We denote by  $\operatorname{Irr}(\mathcal{A}_{\hbar})$  the set of equivalence classes of irreducible algebraic representations of the quantized algebra. We will see that  $\operatorname{Irr}(\mathcal{A}_{\hbar})$  is a set "equivalent" to a torus.

The group  $\Gamma$  naturally acts on a quantized algebra  $\mathcal{A}_{\hbar}$  and hence on the set  $\operatorname{Irr}(\mathcal{A}_{\hbar})$ . Let  $\hbar = \frac{M}{N}$  with  $\gcd(M, N) = 1$ . The following holds:

Theorem 0.3 (Canonical equivariant representation) There exists a unique (up to isomorphism) N-dimensional irreducible representation  $(\pi_{\hbar}, \mathcal{H}_{\hbar})$  of  $\mathcal{A}_{\hbar}$  for which its equivalence class is fixed by  $\Gamma$ .

This means that:

$$\pi_{\scriptscriptstyle\hbar} \simeq \pi_{\scriptscriptstyle\hbar}^B$$

for all  $B \in \Gamma$ .

Since the canonical representation  $(\pi_h, \mathcal{H}_h)$  is irreducible, by Schur's lemma we get the canonical projective representation of  $\Gamma$  compatible with  $\pi_h$ :

Corollary 0.4 (Canonical projective representation) There exists a unique projective representation  $\rho_p: \Gamma \longrightarrow \mathrm{PGL}(\mathcal{H}_{\hbar})$  s.t.:

$$\rho_{\mathbf{p}}(B)^{-1}\pi_{\hbar}(f)\rho_{\mathbf{p}}(B) = \pi_{\hbar}(f \circ B)$$

for all  $f \in \mathcal{A}$  and  $B \in \Gamma$ .

**Remark.** Corollary 0.4 is an improvement to the known constructions (cf. [HB, Me, KR1]) which has the group  $\Gamma_{\Theta} := \{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \mid ab = cd = 0 \ (2) \}$  as the group of quantizable elements.

Using a result of Coxeter-Moser [CM] about the structure of the group  $\Gamma$  we get:

**Theorem 0.5 (Linearization)** The projective representation  $\rho_p$  can be lifted to a linear representation in exactly 12 different ways.

**Remark.** The existence of the linear representation  $\rho_h$  in Theorem 0.5 answers Hannay-Berry's question (cf. [HB, Me]) on the multiplicativity of the map  $\rho_h$ .

**Summary.** For  $\hbar \in \mathbb{Q}$  let  $(\rho_{\hbar}, \pi_{\hbar}, \mathcal{H}_{\hbar})$  be the canonical (projective) equivariant quantization of  $\mathbf{T}$ . We can endow the space  $\mathcal{H}_{\hbar}$  with a canonical unitary structure s.t.  $\pi_{\hbar}$  is a \*-representation and  $\rho_{\hbar}$  is unitary. This "family" of \*-representations of  $\mathcal{A}_{\hbar}$  is by definition a Weyl quantization of the functions on the torus. The above results show the existence of a canonical projective equivariant quantization of the torus, and the existence of a linear equivariant quantization of the torus.

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# 1 Construction

We consider the algebra  $\mathcal{A} := C^{\infty}(\mathbf{T})$  of smooth complex valued function on the torus and the dual lattice  $\Lambda^* := \{\xi \in V^* | \xi(\Lambda) \subset \mathbb{Z}\}$ . Let <, > be the pairing between W and W\*. The map  $\xi \mapsto s(\xi)$  where  $s(\xi)(x) := e^{2\pi i < x, \xi>}$ ,  $x \in \mathbf{T}$  and  $\xi \in \Lambda^*$  defines a canonical isomorphism between  $\Lambda^*$  and the group  $\mathbf{T}^{\vee} := \operatorname{Hom}(\mathbf{T}, \mathbb{C}^*)$  of characters of  $\mathbf{T}$ .

## 1.1 The quantum tori

Fix  $\hbar \in \mathbb{R}$ . The Rieffel's quantum torus (cf. [Ri]) is the non-commutative algebra  $\mathcal{A}_{\hbar}$  defined over  $\mathbb{C}$  by generators  $\{s(\xi), \xi \in \Lambda^*\}$ , and relations:

$$s(\xi + \eta) = e^{\pi i \hbar \omega(\xi, \eta)} s(\xi) s(\eta)$$

for all  $\xi, \eta \in \Lambda^*$ .

Note that the lattice  $\Lambda^*$  serves, using the map  $\xi \mapsto s(\xi)$ , as a basis for the algebra  $\mathcal{A}_{\hbar}$ . This induces an identification of vector spaces  $\mathcal{A}_{\hbar} \cong \mathcal{A}$  for every  $\hbar$ . We will use this identification in order to view elements of the (commutative) space  $\mathcal{A}$  as members of the (non-commutative) space  $\mathcal{A}_{\hbar}$ .

## 1.2 Weyl quantization

To get a Weyl quantization of  $\mathcal{A}$  we use a specific one-parameter family of representations (see subsection 1.4 below) of the quantum tori. This defines an operator  $\pi_{\hbar}(\xi)$  for every  $\xi \in \Lambda^*$ . We extend the construction to every function  $f \in \mathcal{A}$  using the Fourier theory. Suppose:

$$f = \sum_{\xi \in \Lambda^*} a_\xi \cdot \xi$$

is its Fourier expansion. Then we define its Weyl quantization by:

$$\pi_{{\scriptscriptstyle\hbar}}(f) := \sum_{\xi \in \Lambda^*} a_{\xi} \pi_{{\scriptscriptstyle\hbar}}(\xi).$$

The convergence of the last series is due to the rapid decay of the Fourier coefficients of the function f.

# 1.3 Projective equivariant quantization

The group  $\Gamma = \mathrm{SL}_2(\mathbb{Z})$  acts on  $\Lambda$  preserving  $\omega$ . Hence  $\Gamma$  acts on  $\mathcal{A}_{\hbar}$  and the formula of this action is  $s^B(\xi) := s(B\xi)$ . Given a representation  $(\pi_{\hbar}, \mathcal{H}_{\hbar})$  of  $\mathcal{A}_{\hbar}$ 

and an element  $B \in \Gamma$ , define  $\pi_{\hbar}^{B}(s(\xi)) := \pi_{\hbar}(s^{B^{-1}}(\xi))$ . This formula induces an action of  $\Gamma$  on the set  $Irr(\mathcal{A}_{\hbar})$  of equivalence classes of irreducible algebraic representations of  $\mathcal{A}_{\hbar}$ .

**Lemma 1.1** All irreducible representations of  $A_{\hbar}$  are N-dimensional.

Now, suppose  $(\pi_{\hbar}, \mathcal{A}_{\hbar}, \mathcal{H}_{\hbar})$  is an irreducible representation for which its equivalence class is fixed by the action of  $\Gamma$ . This means that for any  $B \in \Gamma$  we have  $\pi_{\hbar} \cong \pi_{\hbar}^{B}$ , so by definition there exists an operator  $\rho_{\hbar}(B) \in GL(\mathcal{H}_{\hbar})$  such that:

$$\rho_{\scriptscriptstyle h}(B)^{-1}\pi_{\scriptscriptstyle h}(\xi)\rho_{\scriptscriptstyle h}(B)=\pi_{\scriptscriptstyle h}(B\xi)$$

for all  $\xi \in \Lambda^*$ . This implies the Egorov identity (0.2.1) for any function . Now, since  $(\pi_{\hbar}, \mathcal{H}_{\hbar})$  is an irreducible representation then by Schur's lemma for every  $B \in \Gamma$  the operator  $\rho_{\hbar}(B)$  is uniquely defined up to a scalar. This implies that  $(\rho_{\hbar}, \mathcal{H}_{\hbar})$  is a projective representation of  $\Gamma$ .

## 1.4 The canonical equivariant quantization

In what follows we consider only the case  $\hbar \in \mathbb{Q}$ . We write  $\hbar$  in the form  $\hbar = \frac{M}{N}$  with  $\gcd(M, N) = 1$ .

**Proposition 1.2** There exists a unique  $\pi_{\hbar} \in \operatorname{Irr}(\mathcal{A}_{\hbar})$  which is a fixed point for the action of  $\Gamma$ .

# 1.5 Unitary structure

Note that  $\mathcal{A}_{\hbar}$  becomes a \*- algebra using the formula  $s(\xi)^* := s(-\xi)$ . Let  $(\pi_{\hbar}, \mathcal{H}_{\hbar})$  be the canonical representation of  $\mathcal{A}_{\hbar}$ .

**Remark 1.3** There exists a canonical (unique up to scalar) unitary structure on  $\mathcal{H}_{\hbar}$  for which  $\pi_{\hbar}$  is a \*-representation.

#### 1.6 Realization

Choosing a symplectic basis for  $\Lambda^*$  we get the identifications  $\Lambda^* \subseteq \mathbb{Z} \oplus \mathbb{Z}$  and  $\Gamma = \mathrm{SL}_2(\mathbb{Z})$ . We will consider the realization on the Hilbert space:

$$\mathcal{H} := L^2(\mathbb{Z}/N\mathbb{Z}).$$

#### 1.6.1 Formula for $\pi$

The representation  $\pi$  is given by:

$$[\pi(m, n)f](x) = \alpha(m, n)\psi(nx)f(x + m),$$

where  $\alpha(\mathbf{m}, \mathbf{n}) := (-1)^{M(\mathbf{m}+\mathbf{n})} e^{\pi i \hbar \mathbf{m} \mathbf{n}}$  and  $\psi(t)$  denote the additive character  $\psi(t) := e^{2\pi i \hbar t}$  on  $\mathbb{Z}/N\mathbb{Z}$ .

#### 1.6.2 formula for $\rho$

The projective representation  $\rho$  is described by the following formulas:

$$\left[\rho\left(\begin{array}{cc} 1 & 1 \\ 0 & 1 \end{array}\right)f\right](x) \ = \ \mathsf{Q}(x)f(x),$$

where  $Q(x) := (-1)^{\varepsilon x} e^{\pi i \hbar x^2}$ , with  $\varepsilon := MN \pmod{2}$ , and:

$$\left[\rho \left(\begin{array}{cc} 0 & -1 \\ 1 & 0 \end{array}\right) f\right](x) = \widehat{f}(x),$$

where  $\hat{f}$  denote the Fourier transform:

$$\widehat{\mathbf{f}}(\mathbf{x}) := \frac{1}{\sqrt{N}} \sum_{\mathbf{y} \in \mathbb{Z}/NZ} \mathbf{f}(\mathbf{y}) \psi(\mathbf{y}\mathbf{x}).$$

# 2 Proofs

### 2.1 Proof of Lemma 1.1

Suppose  $(\pi_{\hbar}, \mathcal{H}_{\hbar})$  is an irreducible representation of  $\mathcal{A}_{\hbar}$ .

Step 1. First we show that  $\mathcal{H}_{\hbar}$  is finite dimensional.  $\mathcal{A}_{\hbar}$  is a finite module over  $Z(\mathcal{A}_{\hbar}) = \{s(N\xi), \xi \in \Lambda^*\}$  which is contained in the center of  $\mathcal{A}_{\hbar}$ . Because  $\mathcal{H}_{\hbar}$  has at most countable dimension (as a quotient space of  $\mathcal{A}_{\hbar}$ ) and  $\mathbb{C}$  is uncountable then by Kaplansky's trick (cf. [MR])  $Z(\mathcal{A}_{\hbar})$  acts on  $\mathcal{H}_{\hbar}$  by scalars. Hence dim  $\mathcal{H}_{\hbar} < \infty$ .

**Step 2.** We show that  $\mathcal{H}_{\hbar}$  is N-dimensional. Choose a basis  $(e_1, e_2)$  of  $\Lambda^*$  s.t.  $\omega(e_1, e_2) = 1$ . Suppose  $\lambda \neq 0$  is an eigenvalue of  $\pi_{\hbar}(e_1)$  and denote by  $\mathcal{H}_{\lambda}$  the corresponding eigenspace. We have the following commutation relation

 $\pi_{\hbar}(e_1)\pi_{\hbar}(e_2) = \gamma \pi_{\hbar}(e_2)\pi_{\hbar}(e_1)$  where  $\gamma := e^{-2\pi i \frac{M}{N}}$ . Hence  $\pi_{\hbar}(e_2) : \mathcal{H}_{\gamma^j \lambda} \longrightarrow \mathcal{H}_{\gamma^{j+1}\lambda}$ , and because  $\gcd(M,N) = 1$  then  $\mathcal{H}_{\gamma^i \lambda} \neq \mathcal{H}_{\gamma^j \lambda}$  for  $0 \leq i \neq j \leq N-1$ . Now, let  $v \in \mathcal{H}_{\lambda}$  and recall that  $\pi_{\hbar}(e_2)^N = \pi_{\hbar}(Ne_2)$  is a scalar operator. Then the space  $\operatorname{span}\{v,\pi_{\hbar}(e_2)v,\ldots,\pi_{\hbar}(e_2)^{N-1}v\}$  is N-dimensional  $\mathcal{A}_{\hbar}$ —invariant subspace hence it equals  $\mathcal{H}_{\hbar}$ .

## 2.2 Proof of Proposition 1.2

Let us show the existence of a unique fixed point for the action of  $\Gamma$  on  $Irr(\mathcal{A}_{\hbar})$ .

Suppose  $(\pi_{\hbar}, \mathcal{H}_{\hbar})$  is an irreducible representation of  $\mathcal{A}_{\hbar}$ . By Schur's lemma for every  $\xi \in \Lambda^*$  the operator  $\pi_{\hbar}(N\xi)$  is a scalar operator, i.e.,  $\pi_{\hbar}(N\xi) = q_{\pi_{\hbar}}(\xi) \cdot \mathbf{I}$ . We have  $\pi_{\hbar}(0) = \mathbf{I}$  and hence  $q_{\pi_{\hbar}}(\xi) \neq 0$  for all  $\xi \in \Lambda^*$ . Thus to any irreducible representation we have attached a scalar function  $q_{\pi_{\hbar}}: \Lambda^* \longrightarrow \mathbb{C}^*$ . Consider the set  $Q_{\hbar}$  of twisted characters of  $\Lambda^*$ :

$$Q_{\mathbf{h}} := \{q : \Lambda^* \longrightarrow \mathbb{C}^*, \ q(\xi + \eta) = (-1)^{MNw(\xi, \eta)} q(\xi) q(\eta) \}.$$

The group  $\Gamma$  acts naturally on this space by  $q^B(\xi) := q(B^{-1}\xi)$ . It is easy to see that we have defined a map  $\mathbf{q} : \operatorname{Irr}(\mathcal{A}_{\hbar}) \longrightarrow Q_{\hbar}$  given by  $\pi_{\hbar} \mapsto q_{\pi_{\hbar}}$  and it is obvious that this map is compatible with the action of  $\Gamma$ . We use the space of twisted characters in order to give a description for the set  $\operatorname{Irr}(\mathcal{A}_{\hbar})$ :

**Lemma 2.1** The map  $\pi_h \mapsto q_{\pi_h}$  is a  $\Gamma$ -equivariant bijection:

$$\mathbf{q}: \operatorname{Irr}(\mathcal{A}_{\hbar}) \longrightarrow Q_{\hbar}.$$

Now, Proposition 1.2 follows from the following claim:

**Claim 2.2** There exists a unique  $q_o \in Q_h$  which is a fixed point for the action of  $\Gamma$ .

**Proof of Lemma 2.1. Step 1.** The map  $\mathbf{q}$  is surjective. Denote by  $\mathbb{T} := \operatorname{Hom}(\Lambda^*, \mathbb{C}^*)$  the group of complex characters of  $\Lambda^*$ . We define an action of  $\mathbb{T}$  on  $\operatorname{Irr}(\mathcal{A}_{\hbar})$  and on  $Q_{\hbar}$  by  $\pi_{\hbar} \mapsto \chi \pi_{\hbar}$  and  $q \mapsto \chi^N q$ , where  $\chi \in \mathbb{T}$ ,  $\pi_{\hbar} \in \operatorname{Irr}(\mathcal{A}_{\hbar})$  and  $q \in Q_{\hbar}$ . The map  $\mathbf{q}$  is clearly a  $\mathbb{T}$ -equivariant map with respect to these actions. Since  $\mathbf{q}$  is  $\mathbb{T}$ -equivariant, it is enough to show that the action of  $\mathbb{T}$  on  $Q_{\hbar}$  is transitive. Suppose  $q_1, q_2 \in Q_{\hbar}$ . By definition there exists a character  $\chi_1 \in \mathbb{T}$  for which  $\chi_1 q_1 = q_2$ . Let  $\chi$  be one of the N's roots of  $\chi_1$  then  $\chi^N q_1 = q_2$ .

Step 2. The map  $\mathbf{q}$  is one to one. Suppose  $(\pi_{\hbar}, \mathcal{H}_{\hbar})$  is an irreducible representation of  $\mathcal{A}_{\hbar}$ . It is easy to deduce from the proof of Lemma 1.1 (Step 2)

that for  $\xi \notin N\Lambda^*$  we have  $\operatorname{tr}(\pi_h(\xi)) = 0$ . But we know from character theory that an isomorphism class of a finite dimensional irreducible representation of an algebra is recovered from its character. This completes the proof of Lemma 2.1.

**Proof of Claim 2.2.** Uniqueness. Fix  $q \in Q_h$ . The map  $\chi \mapsto \chi q$  give a bijection of  $\mathbb{T}$  with  $Q_h$ . But the trivial character  $\mathbf{1} \in \mathbb{T}$  is the unique fixed point for the action of  $\Gamma$  on  $\mathbb{T}$ .

Existence. Choose a basis  $(e_1, e_2)$  of  $\Lambda^*$  s.t.  $\omega(e_1, e_2) = 1$ . This allows to identify  $\Lambda^*$  with  $\mathbb{Z} \oplus \mathbb{Z}$ . It is easy to see that the function:

$$q_{o}(m,n) = (-1)^{MN(mn+m+n)}$$

is a twisted character which is fixed by  $\Gamma$ . This completes the proof of Claim 2.2 and of Proposition 1.2.

### 2.3 Proof of Theorem 0.5

The theorem follows from the following proposition:

**Proposition 2.3** Fix a projective representation  $\rho_p : \Gamma \longrightarrow GL(\mathcal{H}_{\hbar})$ . Then it can be lifted to a linear representation in exactly 12 ways.

*Proof.* Existence. We want to find constants c(B) for every  $B \in \Gamma$  s.t.  $\rho_h := c(\cdot)\rho_p$  is a linear representation of  $\Gamma$ . This is possible to carry out due to the following fact:

**Lemma 2.4 ([CM])** The group  $\Gamma$  is isomorphic to the group generated by three letters S, B and Z subjected to the relations:  $Z^2 = 1$  and  $S^2 = B^3 = Z$ .

Lemma 2.4  $\Rightarrow$  Existence. We need to find constants  $c_Z, c_B, c_S$  so that the operators  $\rho_{\hbar}(Z) := c_Z \rho_{\rm p}(Z), \ \rho_{\hbar}(B) := c_B \rho_{\rm p}(B), \ \rho_{\hbar}(S) := c_S \rho_{\rm p}(S)$  will satisfy the identities:

$$\rho_{h}(Z)^{2} = I, \ \rho_{h}(B)^{3} = \rho_{h}(Z), \ \rho_{h}(S)^{2} = \rho_{h}(Z).$$

This can be done by taking appropriate scalars.

Now, fix one lifting  $\rho_0$ . Then for the collection of operators  $\rho_{\hbar}(B)$  which lifts  $\rho_p$  define a function  $\chi(B)$  by  $\rho_{\hbar}(B) = \chi(B)\rho_0(B)$ . It is obvious that  $\rho_{\hbar}$  is a representation if and only if  $\chi$  is a character. Thus liftings corresponds to characters. By Lemma 2.4 the group of characters  $\Gamma^{\vee} := \operatorname{Hom}(\Gamma, \mathbb{C}^*)$  is isomorphic to  $\mathbb{Z}/12\mathbb{Z}$ .

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